

# The quantum view of the world 5

## Tunneling

Modern physics for engineers

David Miller

# Tunneling



In quantum mechanics

a particle can get to the other side  
of a “hill”

even though it does not have  
enough energy to get over the  
“hill”

This is “tunneling”

# Tunneling



Tunneling is common in modern electronics, e.g.,  
in “flash” memory

which is written and erased by  
tunneling through dielectrics  
in transistors

where there is undesired  
tunneling through the gate  
oxide

# Tunneling in optics

# Tunneling in optics



Shining a flashlight from underwater  
will lead to total internal reflection  
past a certain angle

But the light “tunnels” a short  
distance into the air

# Tunneling in optics



Putting a piece of glass close to the surface  
but not touching it

can allow light to “tunnel” from the water into the glass

known as “frustrated total internal reflection”

a well-known classical wave phenomenon

A potential barrier of finite height

Nature of solutions for a finite barrier

# Infinitely thick barrier

Suppose we have a barrier of height

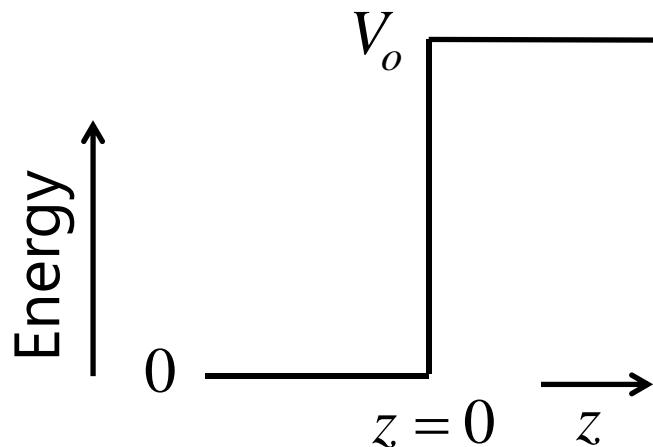
$$V_o$$

with potential 0 to the left of the barrier

A quantum mechanical wave is incident from the left

The energy  $E$  of this wave is positive

i.e.,  $E > 0$



# Infinitely thick barrier

We allow for reflection from the barrier

into the region on the left

Using the general solution on the left

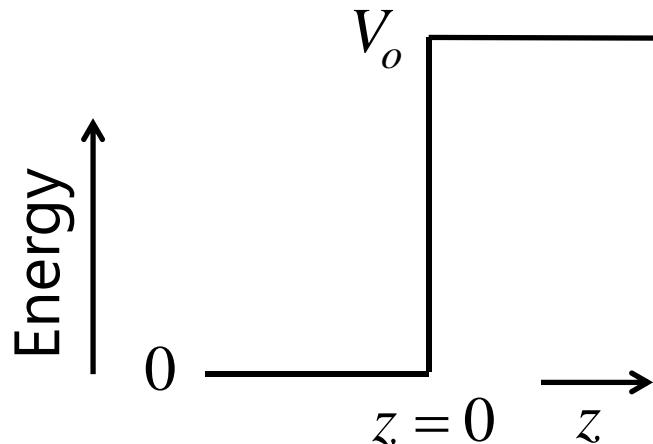
with complex exponential waves

$$\psi_{left}(z) = A \exp(ikz) + B \exp(-ikz)$$

where, as before  $k = \sqrt{2mE / \hbar^2}$

$A \exp(ikz)$  is the incident wave, going right

$B \exp(-ikz)$  is the reflected wave, going left



# Infinitely thick barrier

Presume that  $E < V_o$

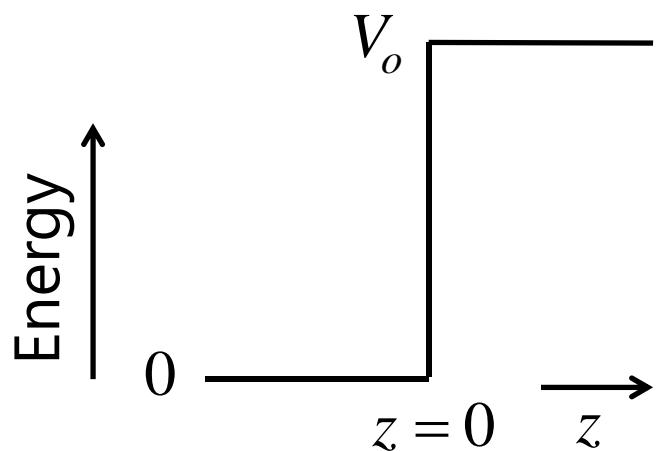
i.e., the incident wave energy is less than the barrier height

Inside the barrier, the wave equation is

$$-\frac{\hbar^2}{2m} \frac{d^2\psi(z)}{dz^2} + V_o\psi(z) = E\psi(z)$$

i.e., mathematically

$$\frac{d^2\psi(z)}{dz^2} = \frac{2m}{\hbar^2} (V_o - E)\psi(z)$$



# Infinitely thick barrier

The general solution of

$$\frac{d^2\psi(z)}{dz^2} = \frac{2m}{\hbar^2} (V_o - E) \psi(z)$$

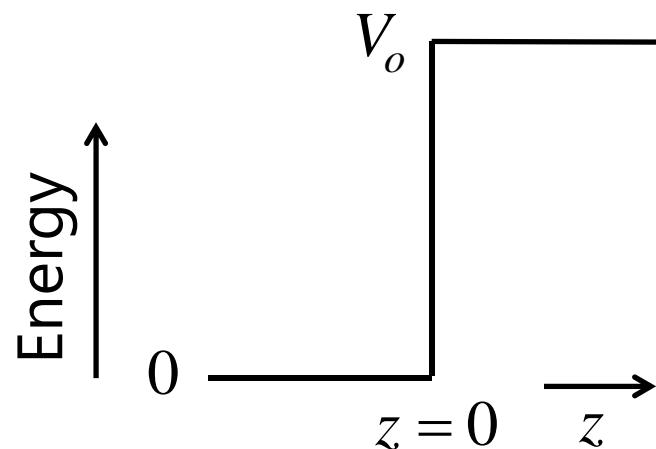
for the wave on the right is

$$\psi_{right}(z) = C \exp(-\kappa z) + D \exp(\kappa z)$$

where  $\kappa = \sqrt{2m(V_o - E) / \hbar^2}$

We presume  $D = 0$

otherwise the wave increases  
exponentially to the right for ever



# Infinitely thick barrier

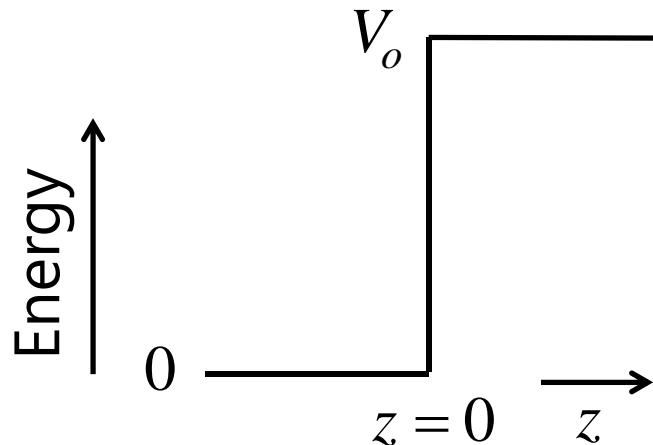
Hence the wave on the right  
inside the barrier, is

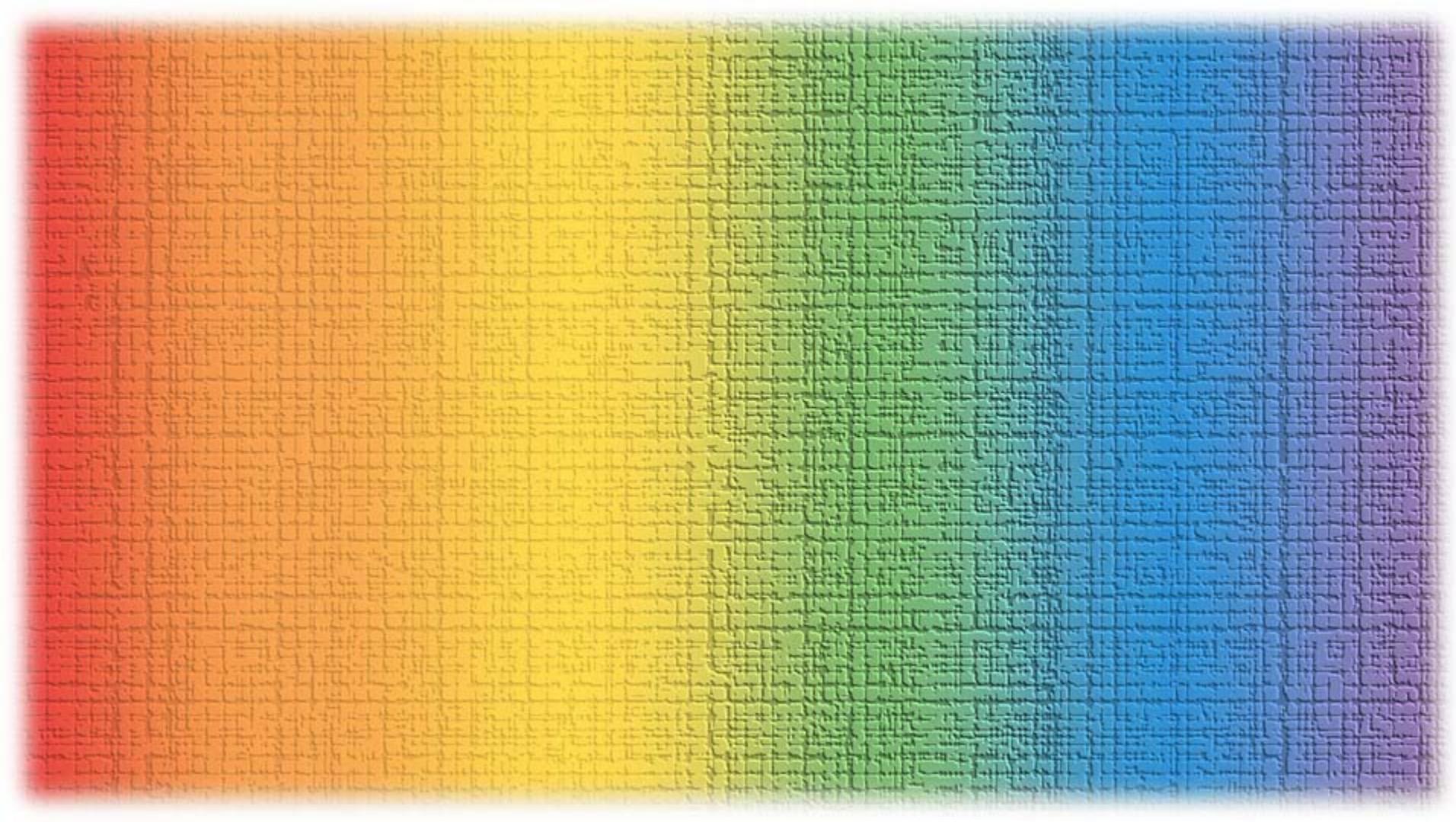
$$\psi_{right}(z) = C \exp(-\kappa z)$$

with  $\kappa = \sqrt{2m(V_o - E) / \hbar^2}$

This solution proposes that the wave  
inside the barrier is not zero

Instead, it falls off exponentially





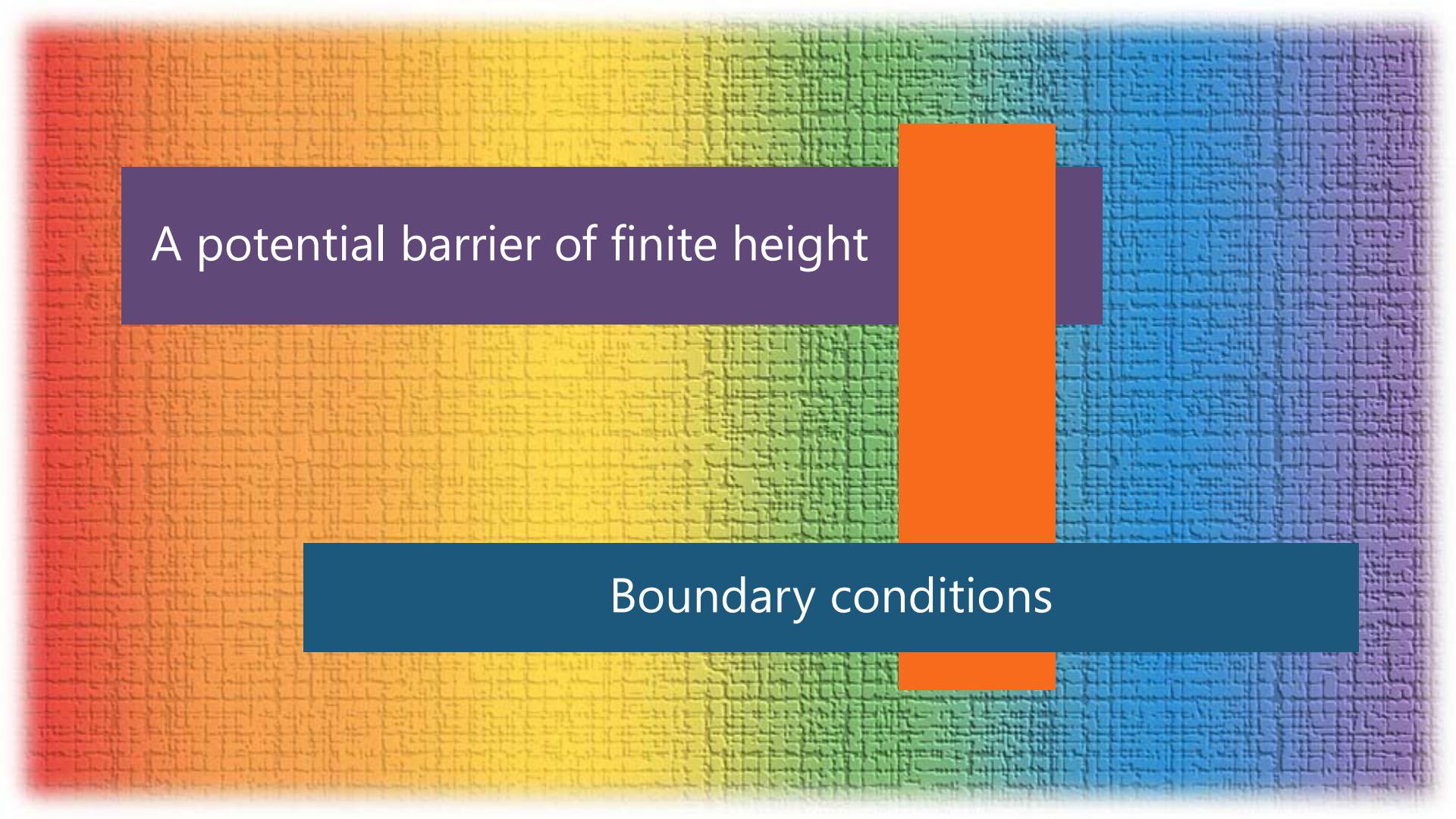


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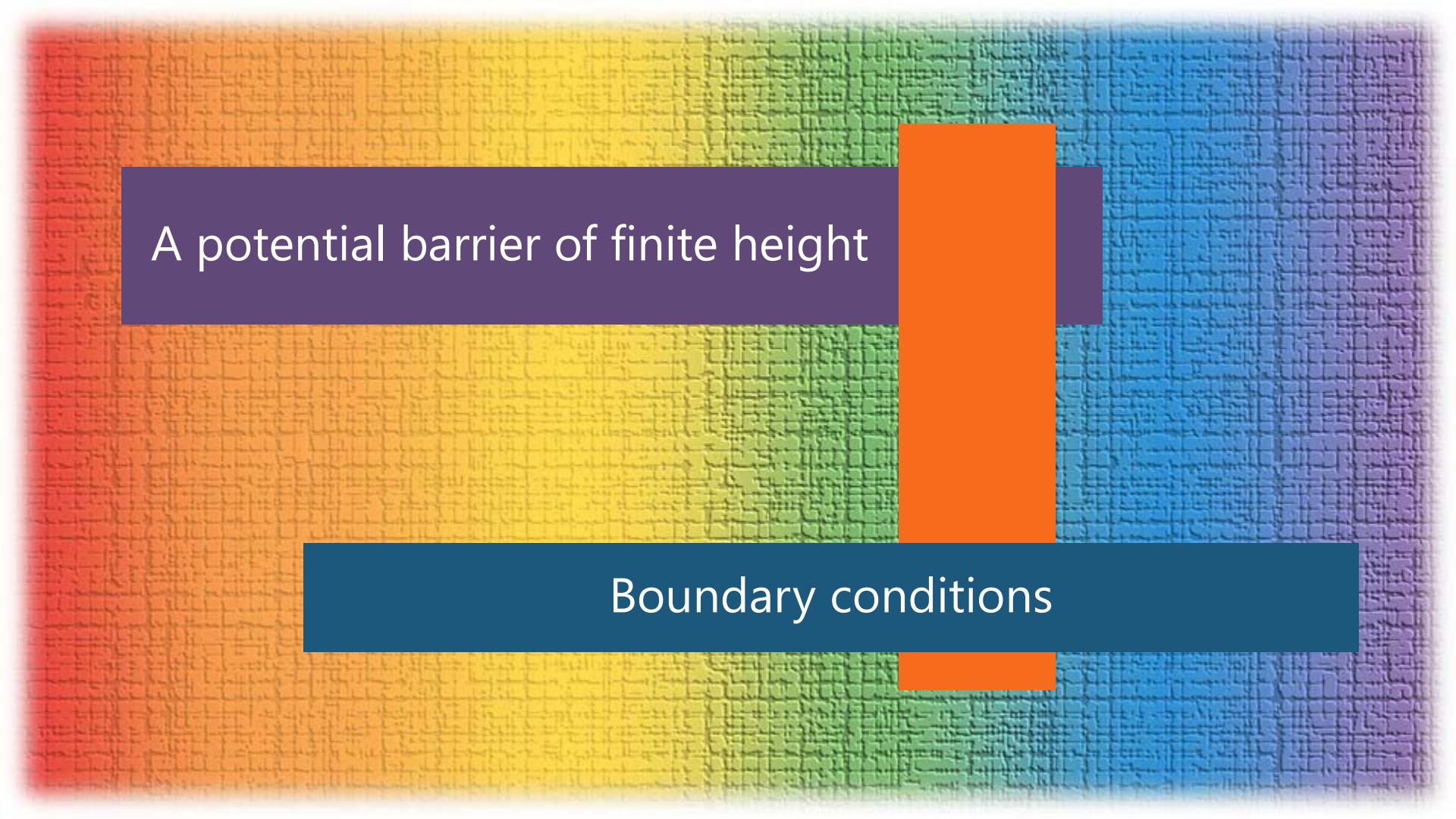
Solving for barriers of finite height

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A potential barrier of finite height



Boundary conditions

# Boundary conditions

For our Schrödinger equation

$$-\frac{\hbar^2}{2m} \frac{d^2\psi(z)}{dz^2} + V(z)\psi(z) = E\psi(z)$$

if we presume that  $E$ ,  $V$  and  $\psi$  are finite  
then  $d^2\psi / dz^2$  must be finite also, so

$d\psi / dz$  must be continuous

If there was a jump in  $d\psi / dz$   
then  $d^2\psi / dz^2$  would be infinite at that point

# Boundary conditions

Also

$d\psi / dz$  must be finite

otherwise  $d^2\psi / dz^2$  could be infinite

being the limit of a difference  
involving infinite quantities

For  $d\psi / dz$  to be finite

$\psi$  must be continuous

# Boundary conditions

Now that we have these two boundary conditions

$\psi$  must be continuous

$d\psi / dz$  must be continuous

we can proceed to solve problems with finite “heights” of boundaries

A potential barrier of finite height

Solutions for a barrier of finite height

# Infinitely thick barrier

Using the boundary conditions

we complete the solution

On the left, we have

On the right we have

Continuity of the wavefunction

at  $z = 0$  gives

Continuity of the wavefunction

derivative at  $z = 0$  gives

i.e.,

$$\psi_{left}(z) = A \exp(ikz) + B \exp(-ikz)$$

$$\psi_{right}(z) = C \exp(-\kappa z)$$

$$A + B = C$$

$$ikA - ikB = -\kappa C$$

$$A - B = \frac{i\kappa}{k} C$$

# Infinitely thick barrier

Adding

$$A + B = C$$

$$A - B = \frac{i\kappa}{k} C$$

gives

$$2A = \left(1 + \frac{i\kappa}{k}\right)C = \left(\frac{k + i\kappa}{k}\right)C$$

Equivalently

$$C = \frac{2k}{k + i\kappa} A = \frac{2k(k - i\kappa)}{k^2 + \kappa^2} A$$

so we have found the amplitude  $C$  of the wave in the barrier in terms of the amplitude  $A$  of the incident wave

# Infinitely thick barrier

Subtracting

$$A + B = C$$

$$A - B = \frac{i\kappa}{k} C$$

gives a similar relation between  $B$  and  $C$   
and we can deduce a relation  
between  $A$  and  $B$

So we can solve the entire problem here  
leaving only one arbitrary overall constant

# Probability densities

Note that we have to take the modulus squared of the entire wavefunction  
so on the left we have

$$|\psi_{left}(z)|^2 = |A \exp(ikz) + B \exp(-ikz)|^2$$

and on the right

$$|\psi_{right}(z)|^2 = |C \exp(-\kappa z)|^2 = |C|^2 \exp(-2\kappa z)$$

Note the probability density decays by  $1/e$   
in a distance  $1/2\kappa$

# Example numbers

For a barrier of height  $V_o = 2$  eV

and an incident electron of energy  $E = 1.5$  eV

On the left  $k \simeq 6.275 \times 10^9 \text{ m}^{-1} \equiv 6.275 \text{ nm}^{-1}$

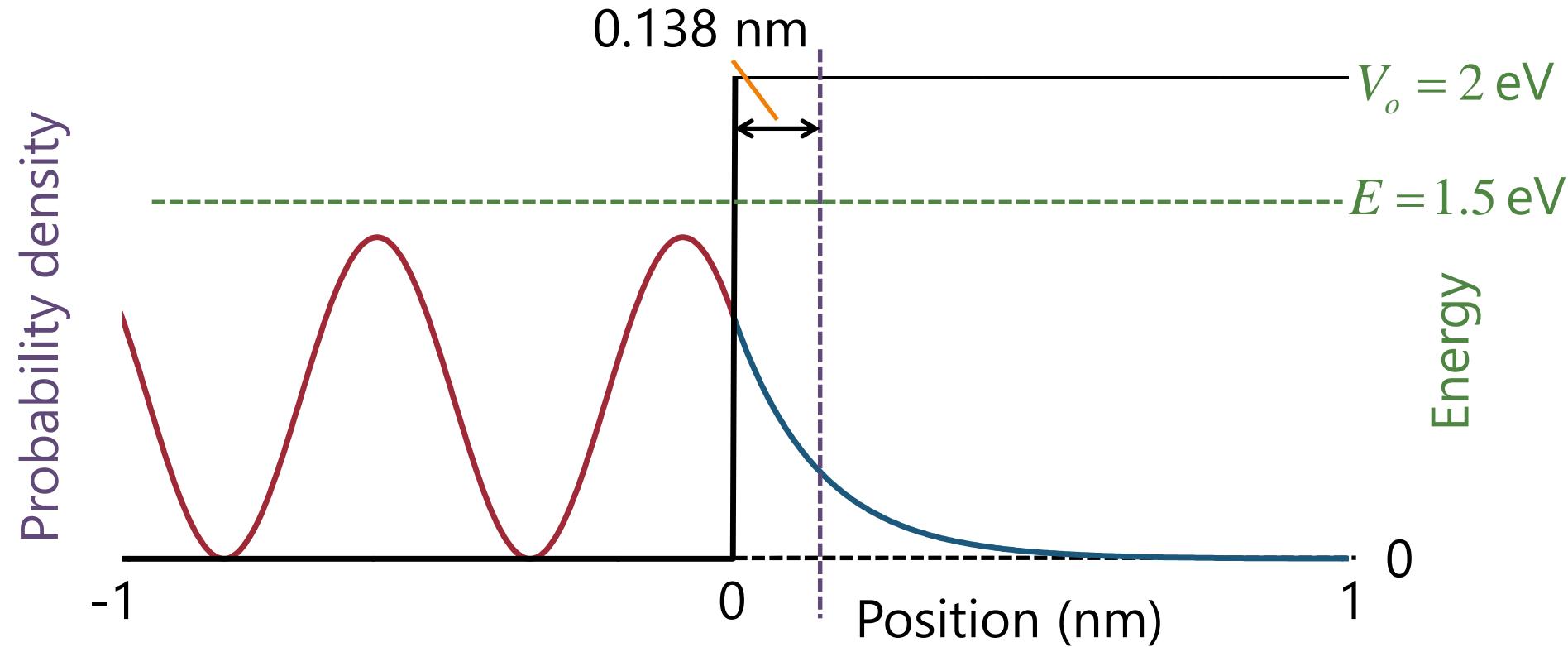
which corresponds to a wavelength of  $\lambda \simeq 1.001 \text{ nm}$

On the right  $\kappa \simeq 3.623 \times 10^9 \text{ m}^{-1} \equiv 3.623 \text{ nm}^{-1}$

so the  $1/e$  decay length of the probability density on the right is  $1/2\kappa \simeq 0.138 \text{ nm}$

# Example probability density calculation

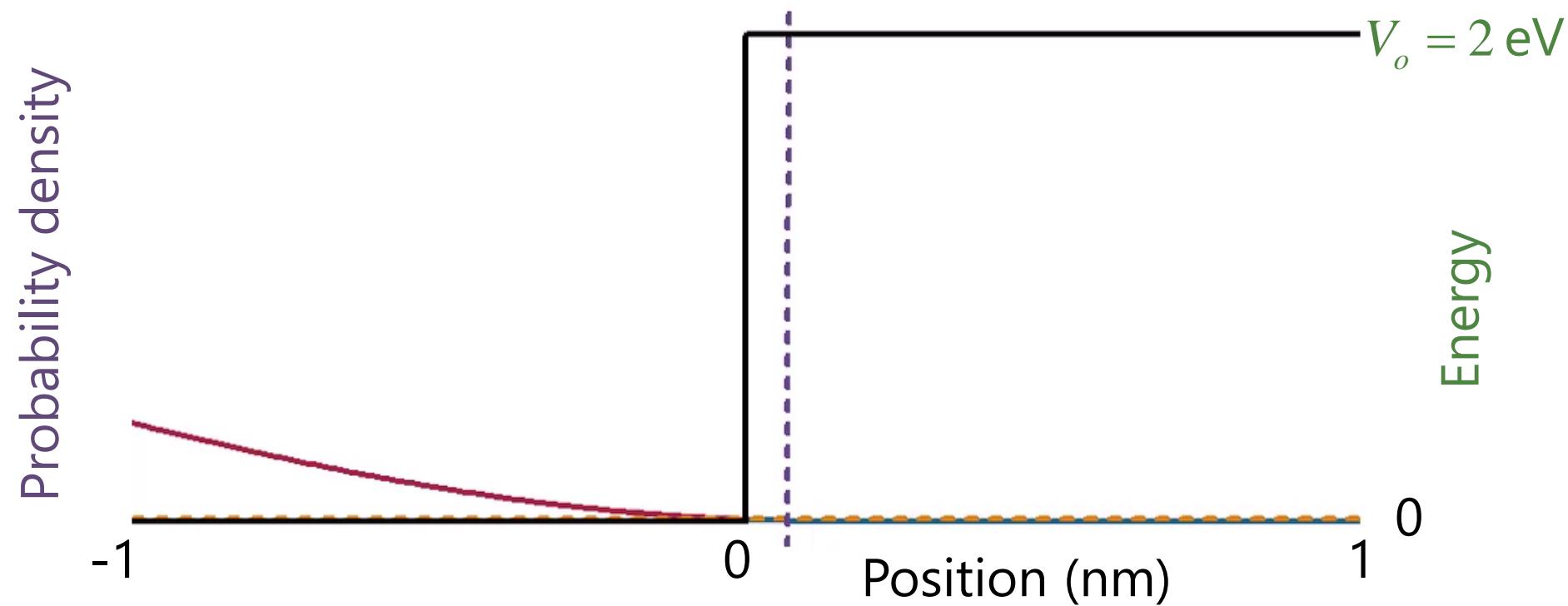
Note the standing wave and the phase change on reflection

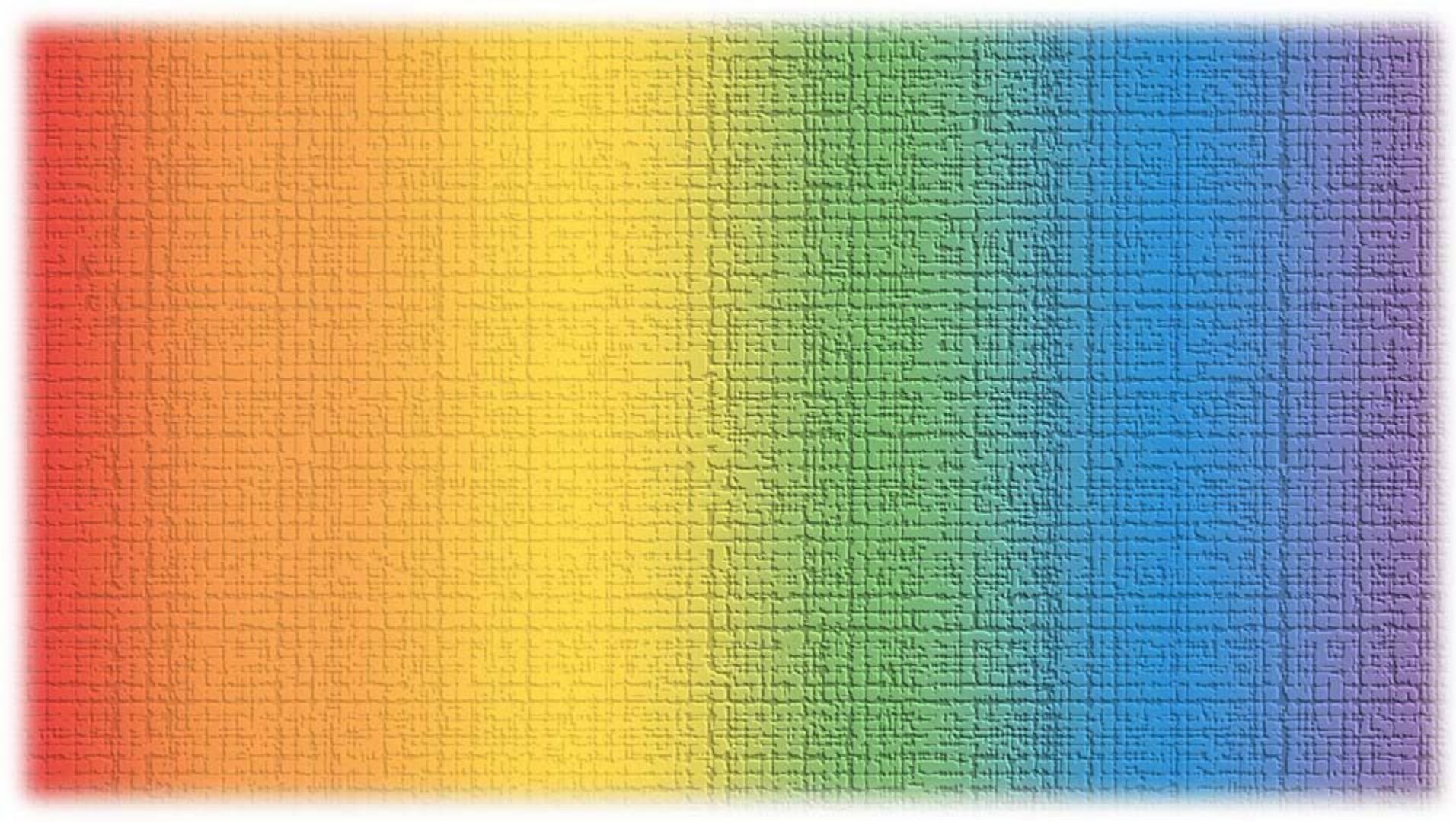


# Example probability density calculation

Energy = 0.01 eV

Penetration depth = 0.069 nm





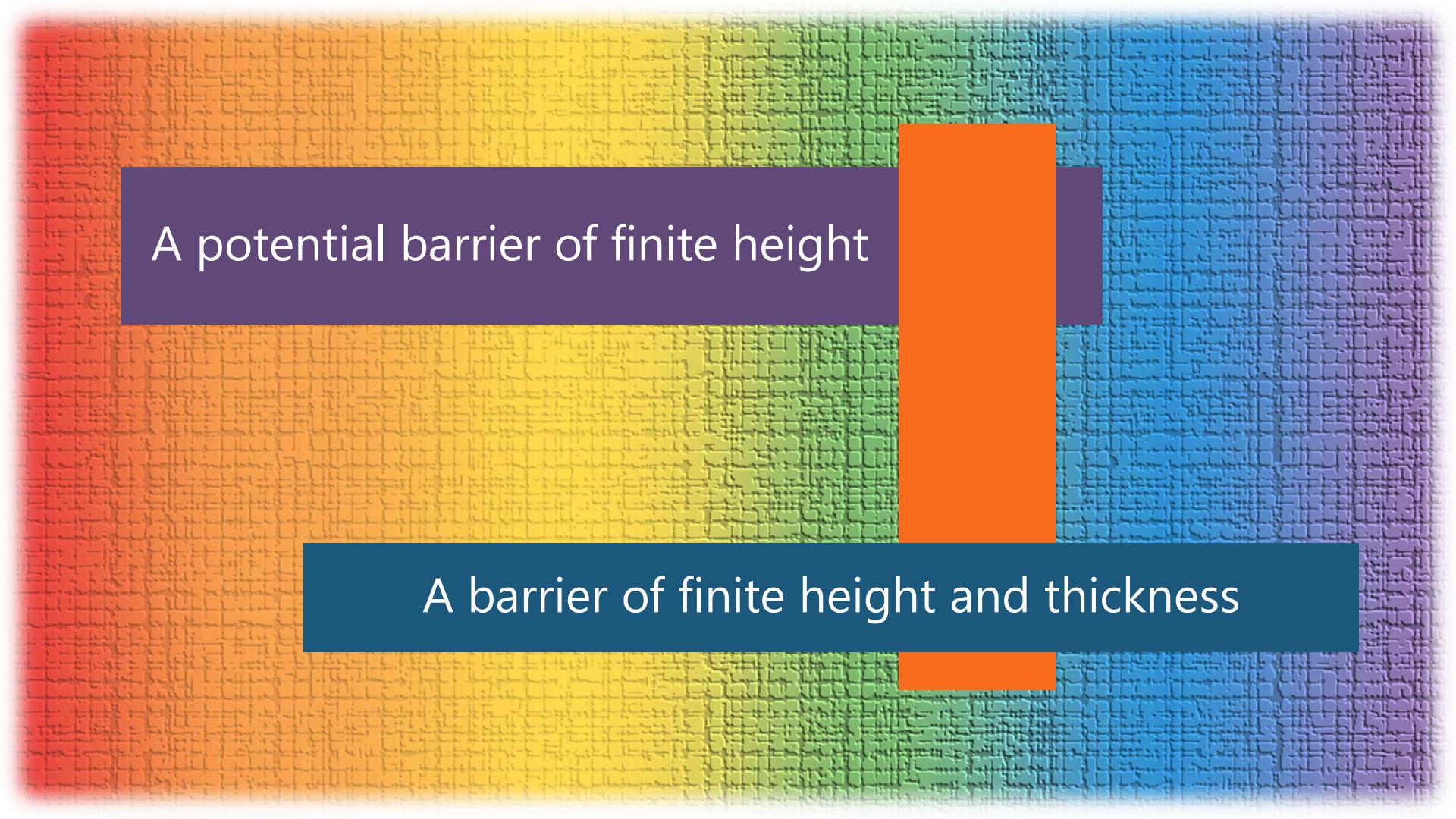


# The quantum view of the world 5

Tunneling through a barrier

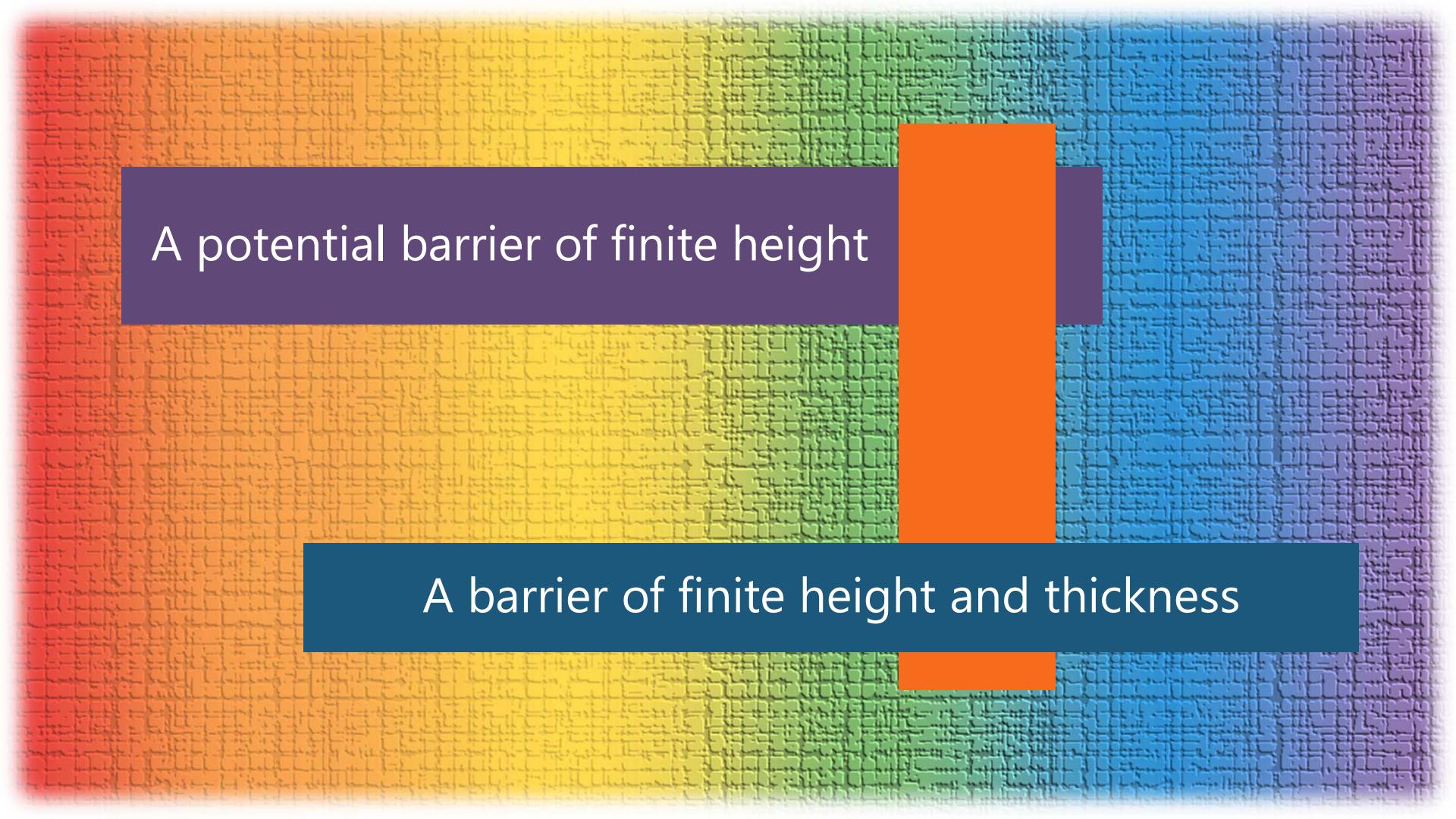
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A potential barrier of finite height

The background is a 2D grid of colored squares, transitioning from red on the left to blue on the right. A central vertical column of squares is colored orange, representing a potential barrier of finite height.



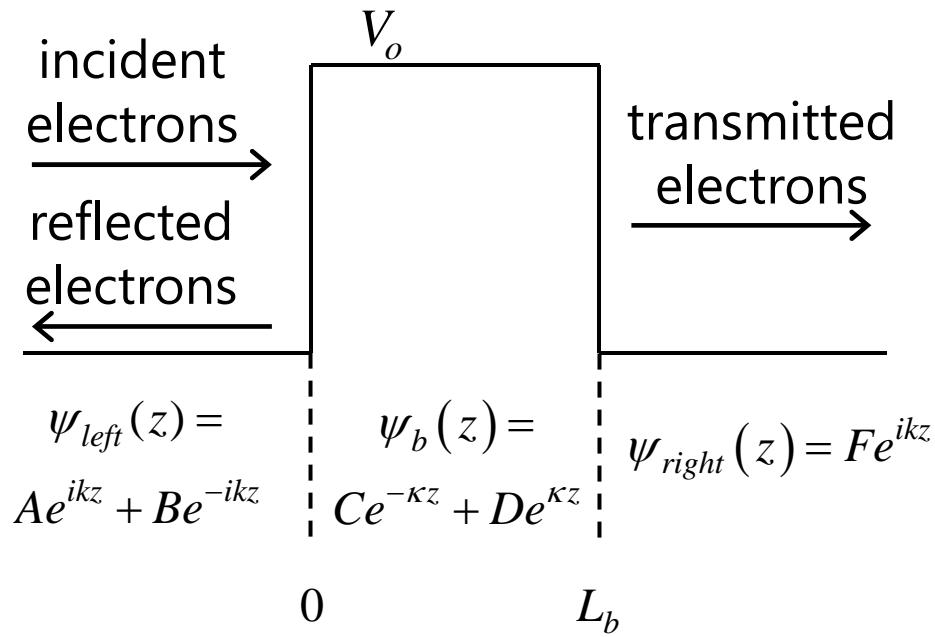
A barrier of finite height and thickness

The background is a 2D grid of colored squares, transitioning from red on the left to blue on the right. A central vertical column of squares is colored orange, representing a barrier of finite height and thickness.

# Tunneling through a barrier

Consider a barrier of finite thickness,  $L_b$   
still with incident electron energy  $E < V_o$   
where  $V_o$  is the barrier height

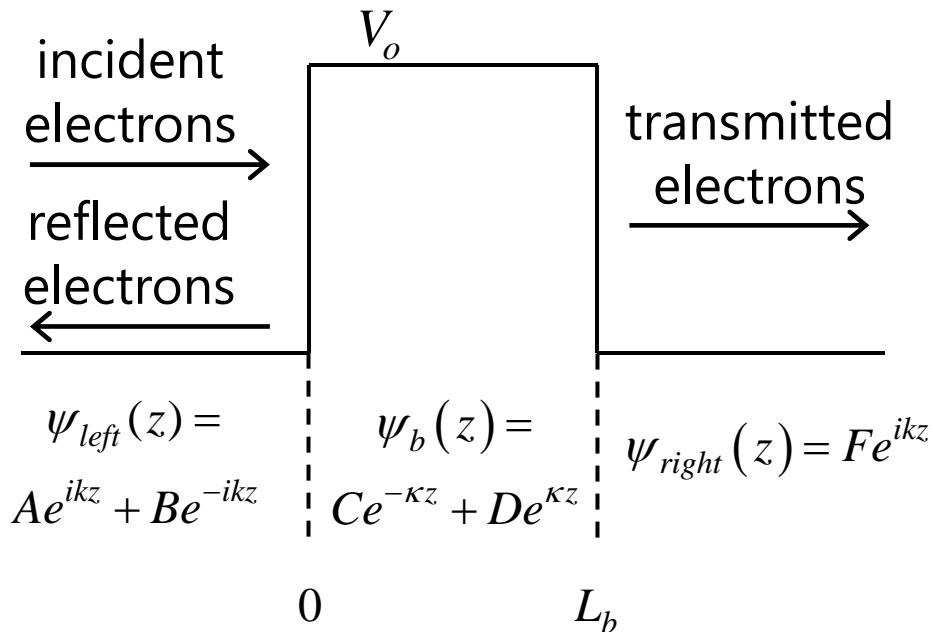
We presume an incident electron wave from the left but none from the right



# Tunneling through a barrier

Now we need to retain both exponentials in the barrier

The “growing” one corresponds to a decaying one from the “reflection” at the right side of the barrier



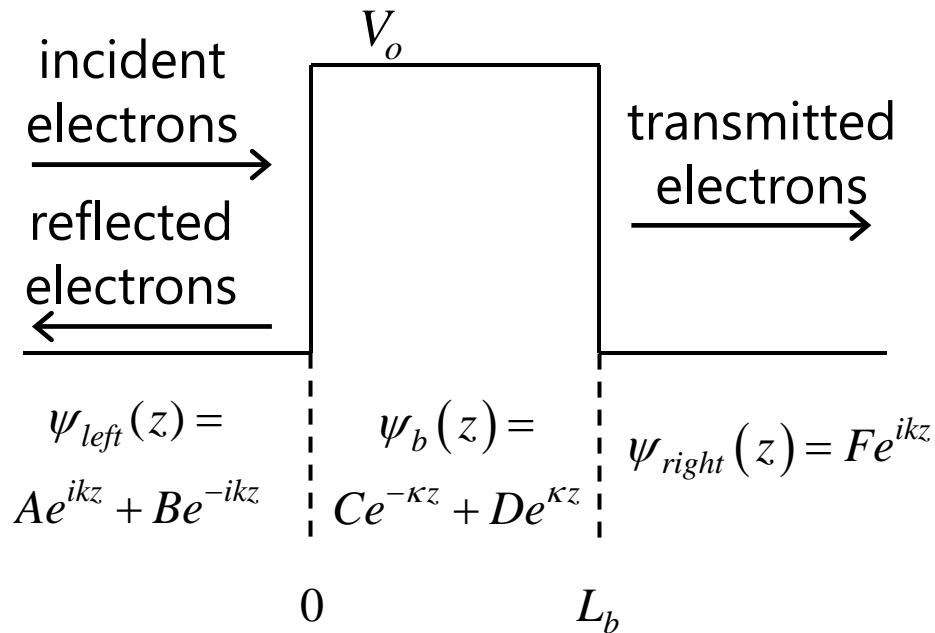
# Tunneling through a barrier

We can solve this starting from the right

Choose an arbitrary amplitude  $F$

Deduce relations between  $C, D$ , and  $F$  using boundary conditions

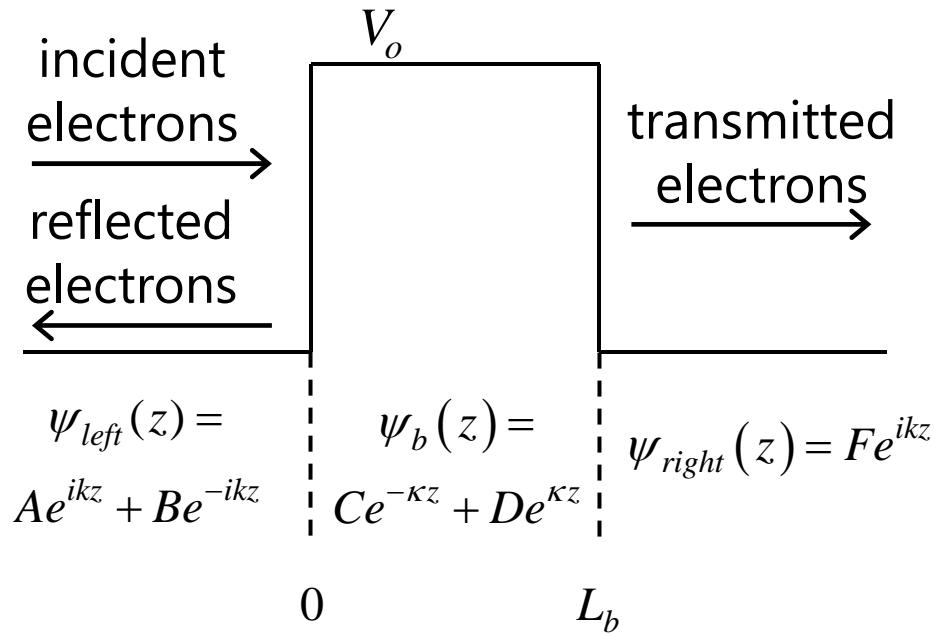
Deduce relations between  $A, B, C$ , and  $D$  using boundary conditions



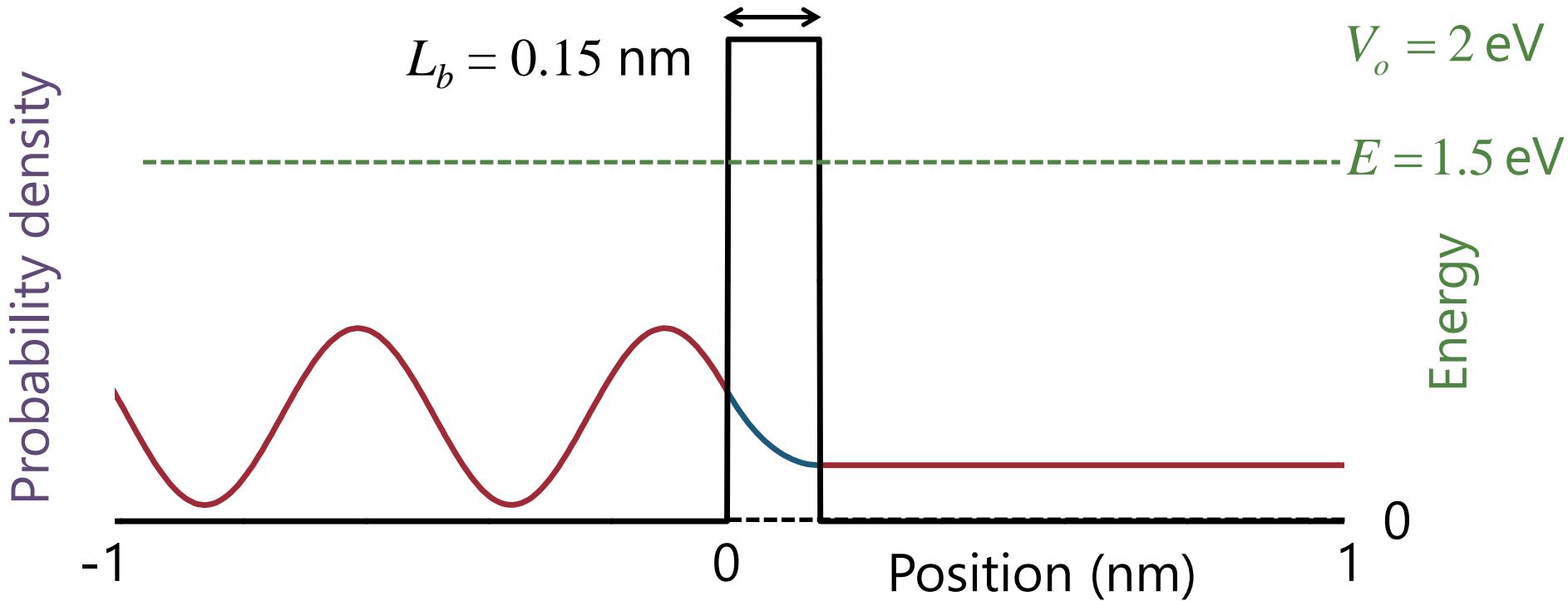
# Tunneling through a barrier

Now the fraction of the incident current formally, in probability density that is transmitted through the barrier will be the ratio  $\frac{|F|^2}{|A|^2}$

We can call this the current “transmission” through the barrier

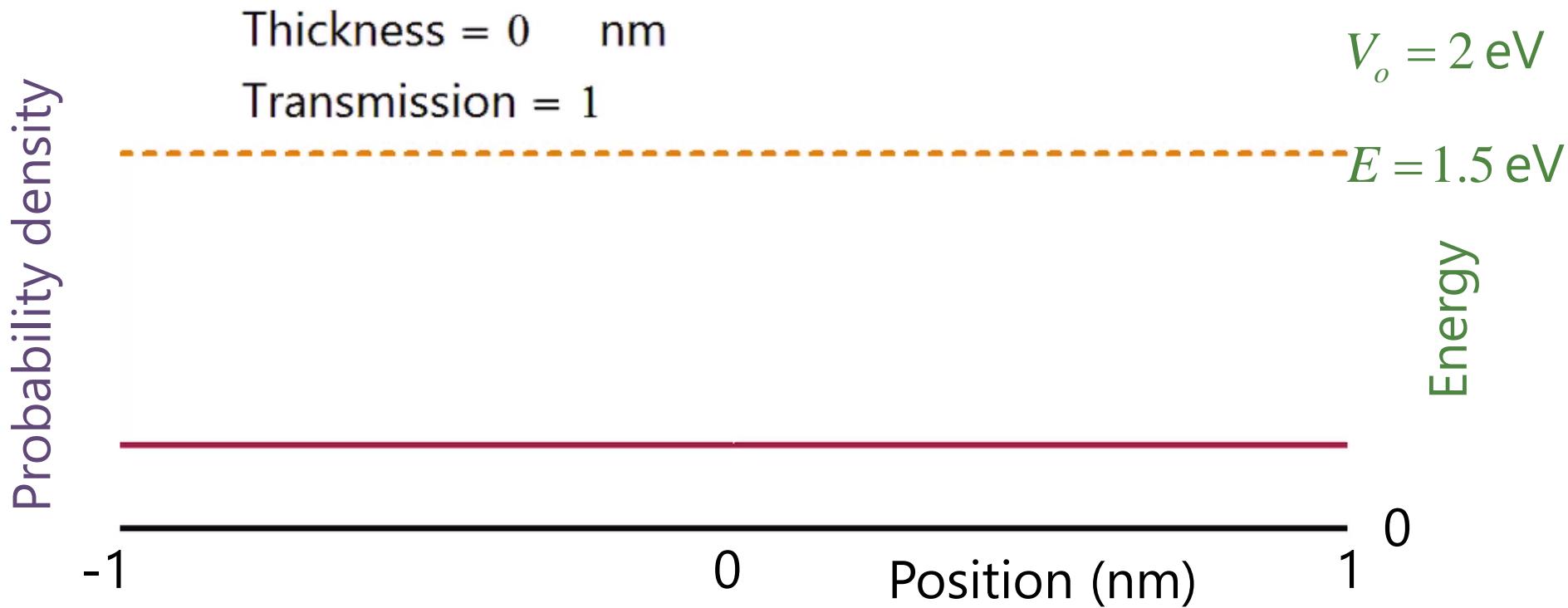


# Tunneling through a barrier



Note the weaker standing wave on the left  
and the transmission to the right

# Tunneling through a barrier



Note the weaker standing wave on the left  
and the transmission to the right

